



Perovskite Photodetectors Based on p-i-n Junction With Epitaxial Electron-Blocking Layers

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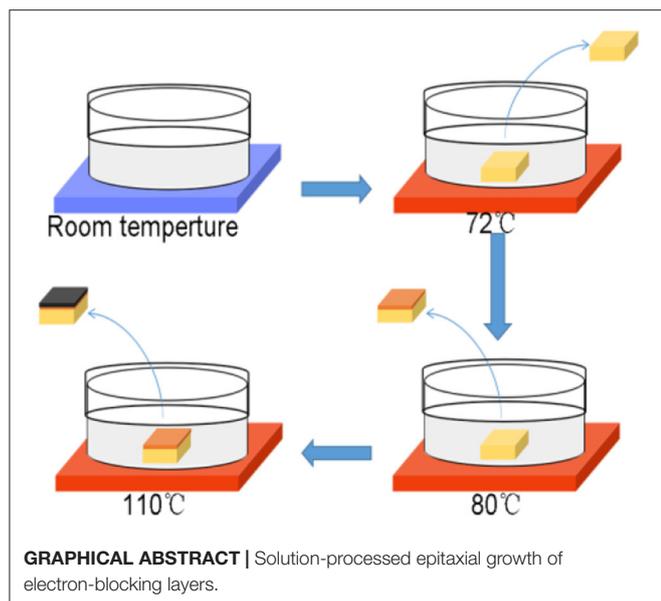
Organic-inorganic hybrid perovskite single crystals (PSCs) have been emerged as remarkable materials for some optoelectronic applications such as solid-state photodetectors, solar cells and light emitting diodes due to their excellent optoelectronic properties. To decrease the dark current, function layers based on spin-coating method are frequently requested for intrinsic PSCs to block the injected current by forming energy barrier. However, the amorphous function layers suffer from small carrier mobility and high traps density, which limit the speed of the photoelectric response of perovskite devices. This work supposes to grow thick MAPbBr₃ and MAPbI₃ mono-crystalline thin films on the surface of intrinsic MAPbBr_{2.5}Cl_{0.5} PSCs substrate by a heteroepitaxial growth technique to act as electron-blocking layers. Meanwhile, C60 and [6,6]-phenyl-C61-butyric acid methyl ester (PCBM) layers are deposited on the opposite surface of substrate PSCs by spin-coating method to block injected holes. This Au-MAPbI₃-MAPbBr₃-MAPbBr_{2.5}Cl_{0.5}PSCs-C60-PCBM-Ag heterostructure can be used as excellent X-ray photodetector (XPD) due to its low dark current density of 6.97×10^{-11} A cm⁻² at -0.5 V bias, high responsivity of 870 mA/W at -100 V bias and X-ray sensitivity as high as 59.7 μC mGy⁻¹ cm⁻² at -50 V bias.

Keywords: photodetectors, self-powered, perovskite single crystals, epitaxial growth, X-ray detection

INTRODUCTION

The last few years have witnessed the intensive investigation on Organic-inorganic hybrid perovskite single crystals (OIHPSCs) by many researchers due to their advantages of high photons absorption coefficients, high charges mobility and adjustable band gaps (Dou et al., 2014). Besides, self-powered photodetectors based on p-n junction and heterostructure can work stably without the applied voltage, which results in their improvement. Moreover, solution-processed optoelectronic devices have shown lower cost compared to those fabricated by vacuum deposition (Dong et al., 2015a; Miao and Zhang, 2019). Especially, OIHPSCs could be grown from solution with bulk size up to centimeter-scale and have shown attractive potential in X-ray detection and gamma ray detection (Wei et al., 2016). In that case, self-powered photodetectors based on solution-processed OIHPSCs tend to have better performance.

The common chemical formulas of OIHP are MAPbX₃ (where MA=CH₃NH₃, X can be Br, Cl or I). The previous studies have demonstrated that diverse solution-processed function layers



deposited on substrate crystals are generally used to fabricate photodetectors because of their transporting of photon-generated charges as well as their blocking of injected current under reverse bias (Wei et al., 2017; Xin et al., 2018; Jariwalad and Hersammc, 2017). To establish p-i-n junctions and decrease the dark current, the function layers based on solution-processed are requested for intrinsic MAPbBr_{2.5}Cl_{0.5} PSCs. This sandwich structure consists of electron transporting layer (ETL), intrinsic PSCs and hole transporting layer (HTL). For example, N₄,N₄₀-bis(4-(6-((3-ethyloxetan-3-yl)methoxy)hexyl)-phenyl)-N₄,N₄₀-diphenylbiphenyl-4,40-diamine (OTPD) and N,N₀-bis(3-methylphenyl)-N,N₀-diphenylbenzidine (poly-TPD), which transport holes, are regarded as p-type semiconductors. And ETLs such as C60, [6,6]-phenyl-C61-butyric acid methyl ester (PCBM) and 2,9-dimethyl-4,7-diphenyl-1,10-phenanthroline (BCP), which transport electrons, are regarded as n-type semiconductors (Wang and Kim, 2017; Wang et al., 2018). In principle, the conduction band energy of intrinsic PSCs should be lower than n-type layer, so that electrons will move to the conduction band of the function layers. And holes will move to the valence band of the inorganic function layers because PSCs' valence band energy should be higher than p-type layer (Sutherland and Sargent, 2016).

However, function layers deposited by spin-coating technique have nanoscale thickness and are easily generating defects, which lead to high traps density and large leakage current (Dong et al., 2015b; Kim et al., 2017). In addition, the function layers have low mobility resulting in difficult transportation of charges. What's more, lattice mismatching is unavoidable between function layers and PSCs substrate. As a result, the devices made by spin-coating method suffer from large inner capacitance due to the accumulation of charges near interface region, which limits the speed of photoelectric response (Xin et al., 2018; Zhao et al., 2014). This work fabricates the photodetector with spin-coated

PCBM/C60 films on MAPbBr_{2.5}Cl_{0.5}PSCs for comparison, the dark current density is $4.24 \times 10^{-8} \text{ A cm}^{-2}$ at low bias voltage, which is not suitable for weak signal detection. Additionally, the weak tolerance to radiation of MAPbBr₃ PSC with spin-coated function layers results in added noise for high-energy ray detection has been reported (Xin et al., 2018). We fabricates the compared device with spin-coated PCBM/C60 films on MAPbBr_{2.5}Cl_{0.5}PSCs Recently, through changing the halide ratio, such as MAPbI_{3-x}Br_x and MAPbBr_{3-x}Cl_x, the band gap of OIHPSCs can be modulated from 1.5 to 2.9 eV (Hu et al., 2016; Fang et al., 2015; Jiang et al., 2019). Among these OIHPSCs, MAPbBr_(x<2.5)Cl_(3-x) PSCs, OIHP single crystals become more n-type with x decreasing. On the other hand, MAPbBr_(x>2.5)Cl_(3-x) and MAPbBr_(3-x)I_x PSCs become more p-type with x increasing as shown in **Figure 1A** (Wang et al., 2014; Jiang et al., 2019).

Therefore, we suppose to employ heteroepitaxial growth technique by solution-processed to form p-i-n junctions. According to the lattice match rule, high quality perovskite layers have to be grown directly on the surface of substrate OIHPSCs in its precursor solution during crystallization process. In this work, we demonstrate that epitaxial perovskite films of high crystallinity could be grown on OIHPSCs. Compared with the spin-coated function layers, epitaxial perovskite films increase the carrier mobility and decrease the trap density. By controlling the time of epitaxial growth in precursor solution, the thickness of perovskite layers can be controlled in micrometer scale, which ensures the devices work stably under high voltage bias. In that case, our perovskite device of centimeter-scale thickness is mainly dedicated to high-energy ray detection such as X-ray detection. Herein, we fabricate the XPD with MAPbBr_{2.5}Cl_{0.5} PSCs as substrate crystals. Then epitaxial MAPbBr₃ mono-crystalline thin films and MAPbI₃ mono-crystalline thin films replace spin-coated organic function layers as p-type, meanwhile PCBM and C60 films act as n-type. The epitaxial MAPbBr₃ layers are used as buffer layers to satisfy the lattice matching condition for epitaxial MAPbI₃ layers and MAPbBr_{2.5}Cl_{0.5} PSCs. This XPD has shown low dark current density of $6.97 \times 10^{-11} \text{ A cm}^{-2}$ at -0.5 V bias, average electrons mobility of $385.9 \text{ cm}^2 \text{ V}^{-1} \text{ s}^{-1}$, high responsivity of 870 mA/W at -100 V bias, average noise current of $1.34 \text{ nA Hz}^{-0.5}$ in darkness under -100 V bias as well as high X-ray sensitivity of $59.7 \mu\text{C mGy}^{-1} \text{ cm}^{-2}$ at -50 V bias.

MATERIALS AND METHODS

Materials: Lead iodide (PbI₂, 99%), lead bromine (PbBr₂, 98%) and Lead chloride (PbCl₂, 99%) were purchased from Sigma Aldrich. Methylamine alcohol solution (30–33 wt. % in ethanol) and Hydroiodic Acid (40%) and Hydrobromic Acid solution (55%) were obtained from Macklin. PCBM and C60 were purchased from Aladdin. All commercial products were used as received.

Device fabrication: The precursor is a N,N-dimethylformamide (DMF) solution which dissolves 1M MABr, PbBr₂ and PbCl₂ with a ratio of 1:0.75:0.25. After totally

reacting and cooling down, the precursor solution were heated in crystallizing dish for 90 h from 50 to 75°C at a rate of 0.3°C h⁻¹ for growth of high-quality MAPbBr_{2.5}Cl_{0.5} PSCs. 1M MABr, PbBr₂ with a ratio of 1:1 are dissolved in a DMF solution for epitaxial growth of MAPbBr₃ mono-crystalline thin films at 80°C. And 1M MAI, PbI₂ with a ratio of 1:1 are dissolved in a γ -GBL solution to grow MAPbI₃ mono-crystalline thin films at 110°C. The programmable heating control system was realized on IKA-RET control-visc.

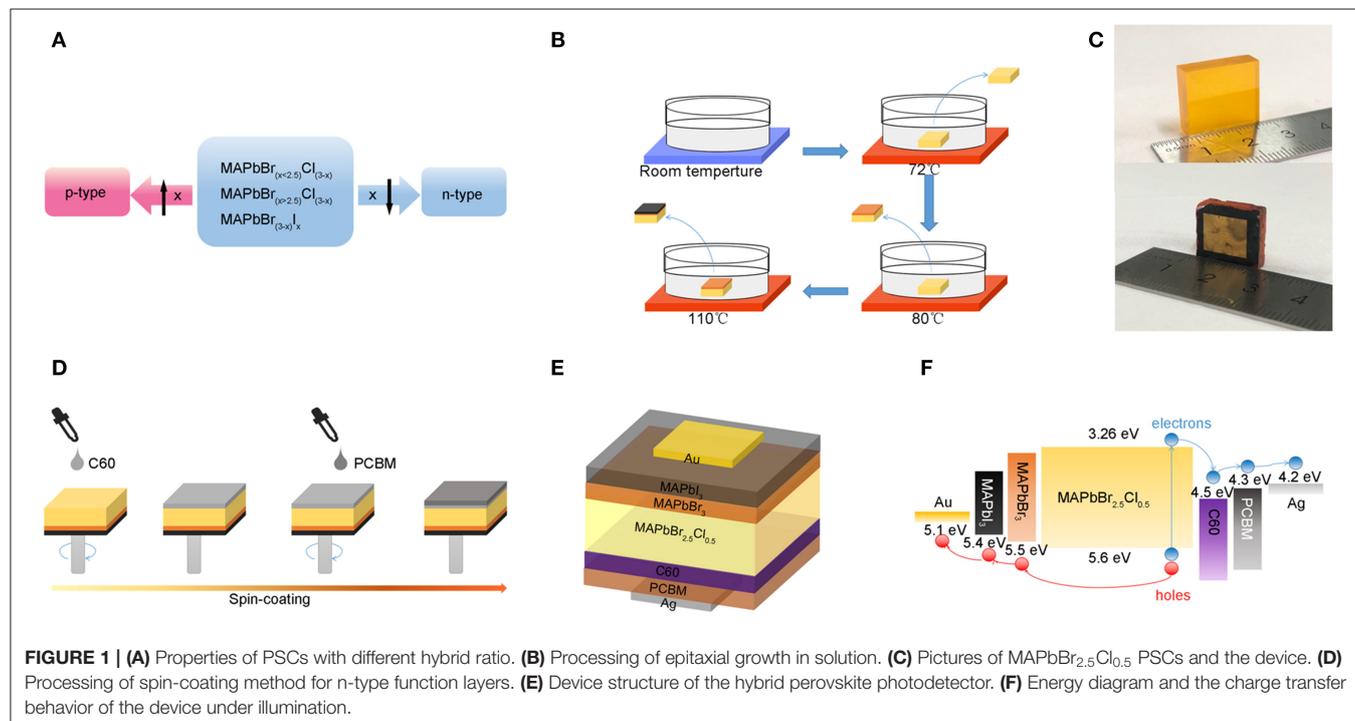
Spin-coating method is used to form C60 and PCBM function films respectively on the bottom surface of MAPbBr_{2.5}Cl_{0.5} PSCs at 2500 r.p.m for 20s. C60, PCBM and silver electrode solution were deposited on the opposite face by thermal evaporation in vacuum. MAPbBr₃ mono-crystalline thin films, MAPbI₃ mono-crystalline thin films and gold electrode were deposited on opposite surfaces of PSCs.

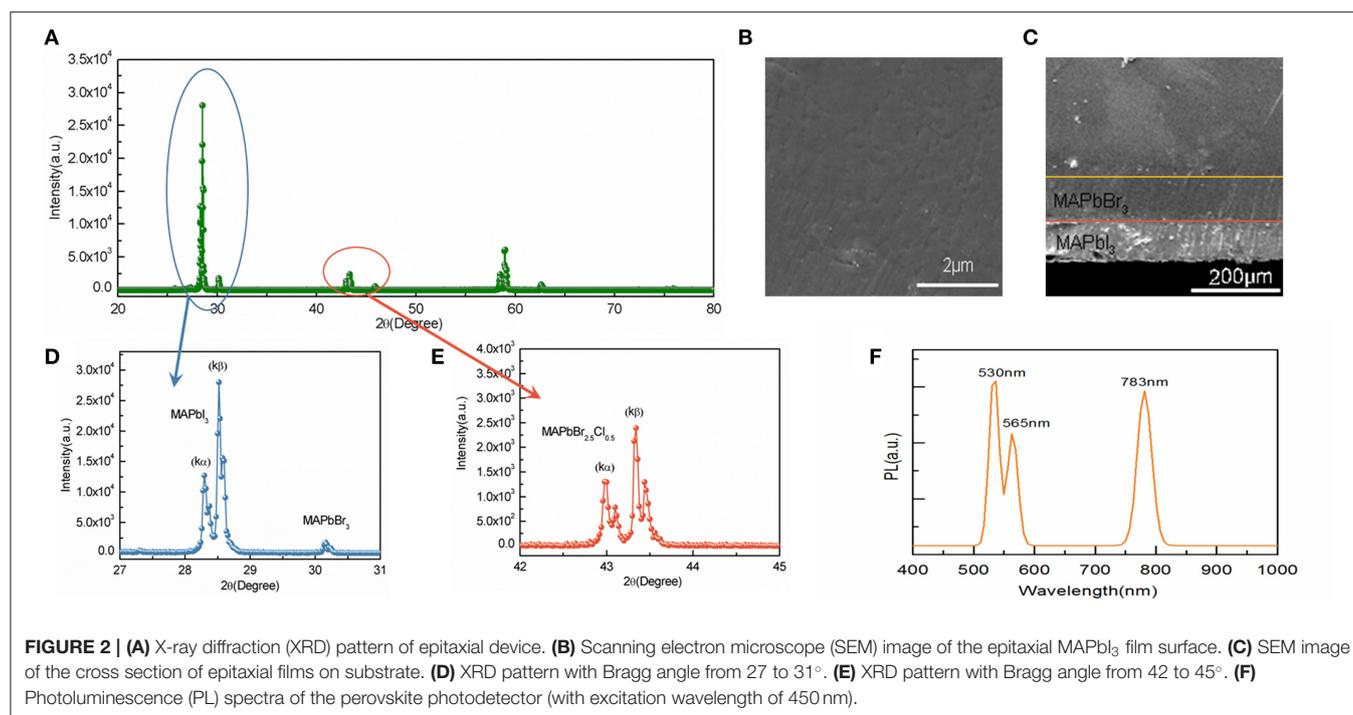
Characterization: SEM images were taken with a Quanta 200 FEI (USA). XRD patterns were obtained by X'TRA (Switzerland). PL spectra were measured by UV-vis spectroscopy (Lab Tech Bluestar, USA).

Device Performance Measurements: Keithley 2410 was used as the voltage source to measure J-V curves. Noise of the device was characterized with a 450 nm LED at 10 Hz and Keithley 2410 as voltage source. The 460 nm light source is obtained from the F5 LED of GEESLED Company. A 355 nm pulsed Nd:YAG laser with 6 ns pulse width at 20 Hz was used as the illumination source and response time was measured using an Agilent oscilloscope with a Keithley 2410 as the voltage source. The X-ray source is obtained from PERLOVE Medical Company and the X-ray dose rate was obtained by a commercial dosimeter (FJ-347A, China).

RESULTS AND DISCUSSION

PSCs used in this work are grown by variable temperature crystallization (VTC) method which has been previously reported. (Stoumpos et al., 2013; Maculan et al., 2015) The fabrication procedure of MAPbBr_{2.5}Cl_{0.5} PSCs and epitaxial perovskite mono-crystalline films is depicted in **Figure 1B**. As illustrated in **Figure 1C**, the size of MAPbBr_{2.5}Cl_{0.5} PSC is approximately 18 mm × 18 mm in width and length, as well as 5 mm in thickness. The MAPbBr_{2.5}Cl_{0.5} PSC is dipped into MAPbBr₃ and MAPbI₃ perovskite precursor solution respectively at optimum growth temperature. Bragg equation is used to calculate the lattice constants of MAPbBr_{2.5}Cl_{0.5}, MAPbBr₃ and MAPbI₃ (Jiang et al., 2019; Ji et al., 2018). Additionally, X-ray diffraction patterns of each single crystal are shown in **Figure S1** for comparison. And the lattice constants of MAPbBr_{2.5}Cl_{0.5}, MAPbBr₃ and MAPbI₃ are 5.87Å, 5.93 Å and 6.23 Å, respectively. Besides, the lattice mismatching rates of MAPbBr_{2.5}Cl_{0.5}-MAPbBr₃ and MAPbBr₃-MAPbI₃ are calculated to be only 1.01 and 4.81%. As a result, MAPbBr₃ and MAPbI₃ mono-crystalline thin films are formed on the surface of the MAPbBr_{2.5}Cl_{0.5} PSCs in sequence. After cutting and polishing process, the epitaxial mono-crystalline thin films exist only on the top of substrate. As shown in **Figure 1D**, C60 and PCBM films, which act as n-type materials, are deposited on the bottom of MAPbBr_{2.5}Cl_{0.5} PSCs via spin-coating method. The gold and silver electrodes are deposited on the top and bottom surfaces of the PSCs by thermal evaporation in vacuum. **Figure 1E** depicts the Au-MAPbI₃-MAPbBr₃-MAPbBr_{2.5}Cl_{0.5}-C60-PCBM-Ag structure of





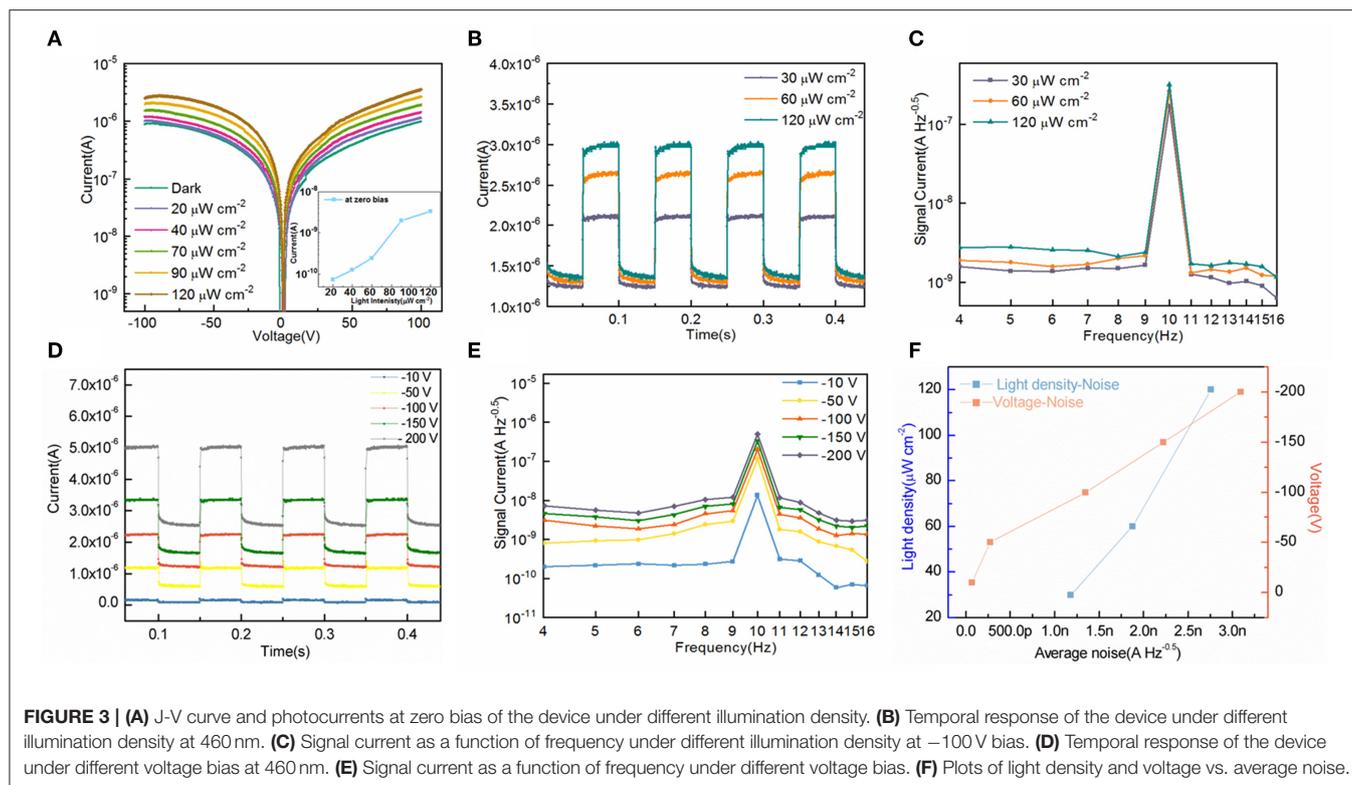
our photodetector. The energy diagram and the charge transfer behavior of this XPD with incident light from silver electrode are shown in **Figure 1F**. These photogenerated electrons and holes are transmitted by ETLs and HTLs and then are collected separately by the electrodes under external electric field. Moreover, C60 and PCBM films can block holes while MAPbI₃ and MAPbBr₃ mono-crystalline thin films block electrons injected by voltage source, which mainly benefit dark current (Bao et al., 2017).

As shown in **Figure 2A**, X-ray diffraction (XRD) is used to characterize the crystallinity of MAPbBr₃ films and MAPbI₃ films on the surface of MAPbBr_{2.5}Cl_{0.5} PSCs. Scanning electron microscope (SEM) characterization of the surface of MAPbI₃ mono-crystalline thin films layer on MAPbBr₃ mono-crystalline thin films is shown in **Figure 2B**. Moreover, **Figure 2C** illustrates the cross-sectional view of the device in which the thickness of each epitaxial crystal layer is clearly visible, with approximately 110 μm of MAPbBr₃ film and 100 μm of MAPbI₃ film. Additionally, Energy Dispersive X-Ray (EDX) spectra of substrate and each layer are illustrated in **Figure S1** to confirm the layers' thickness. Herein, it can be concluded by our experiments that epitaxial mono-crystalline thin films based on lattice matching have considerable thickness, excellent crystallinity (Pellegrino et al., 2015). Additionally, **Figure 2D** illustrates that the epitaxial films deposited on structure have a clear MAPbI₃ peak with full width at half-maximum (FWHM) of 0.105° and this strong peak is separated into kα and kβ peaks, which confirms high crystallization quality. In addition, there is a MAPbBr₃ peak with FWHM of 0.062°. Similarly, as shown in **Figure 2E**, the FWHM of MAPbBr_{2.5}Cl_{0.5} is 0.112°. For comparison,

the XRD patterns of each single crystal are shown in **Figure S2**.

Figure 2F shows the Photoluminescence (PL) spectra (excited by 450 nm) of MAPbBr_{2.5}Cl_{0.5} PSC substrate with epitaxial growth of MAPbBr₃ and MAPbI₃ mono-crystalline thin films. Three emission peaks are obtained at wavelength of 530 nm, 565 nm and 783 nm respectively with a LED excitation at 450 nm. In addition, the PL spectra of each single crystal is provided in **Figure S3**. For comparison, the PL peaks of this heterostructure device obviously become more red-shifted than those of perovskite single crystals. According to previous literatures, the soft and dynamic crystal lattice of halide perovskites are attributed to the low energy barrier for perovskite crystal formation, which easily leads to solid-state ion migrations in the adjacent perovskite lattice (Shewmon et al., 2016; Pan et al., 2018; Zhang et al., 2019). During the process of epitaxial growth in this work, the MAPbBr_{2.5}Cl_{0.5}-MAPbBr₃ and the MAPbBr₃-MAPbI₃ perovskite heterostructures are fabricated respectively where chloride-bromide and bromide-iodide interdiffusion across the perovskite heterojunction occur which leads to the red-shifted of PL peaks.

To characterize the electric properties of this photodetector for visible light, a white light emitting diode with adjustable illumination density is used to radiate gold electrode of the device. Due to the dark current density of photodetectors should be as low as possible in order to recognize weak signal current from noise (Nakamura et al., 2001; Li et al., 2020). The current density-voltage curves of our device under the sweep voltage from -100 V to 100 V at different illumination density are shown in **Figure 3A**. Thanks to the epitaxial mono-crystalline films working as blocking layers, the dark current density of this device



is fairly low of $6.97 \times 10^{-11} \text{ A cm}^{-2}$ at -0.5 V bias, which is one tenth of the $\text{CH}_3\text{NH}_3\text{PbI}_3$ perovskite device reported before only with thick PCBM (50 nm)/C60 (130 nm) layer as the ETL (Lin et al., 2015). What's more, we fabricate Device A with spin-coated PCBM/C60 films on the $\text{MAPbBr}_{2.5}\text{Cl}_{0.5}$ substrate as the ETLs for comparison. The picture and J-V curve of Device A are shown in **Figure S4**. The dark current density of Device A at -0.5 V bias is $4.24 \times 10^{-8} \text{ A cm}^{-2}$, which is three orders of magnitude higher than our photodetector based on epitaxial mono-crystalline films. When the illumination density increases from $20 \mu\text{Wcm}^{-2}$ to $120 \mu\text{Wcm}^{-2}$, the photocurrent density increase from $1.81 \mu\text{A cm}^{-2}$ to $5.12 \mu\text{A cm}^{-2}$ at -50 V bias. Besides, the device could work at zero-bias as shown in the inset due to the in-built electrical field.

Figure 3B illustrates the temporal response of photodetectors based on epitaxial crystal layers with increased density of 460 nm illumination at 10 Hz. In order to get the signal noise at low frequency of the device, the Fast Fourier Transform (FFT) method is used to transfer photocurrent density vs. time into signal current density vs. frequency (Dong et al., 2015b). And we analyze the signal current density ranges from 1 to 16 Hz, as shown in **Figure 3C**. The noise current at low frequency increases approximately from 0.62 to 1.54 nA $\text{Hz}^{-0.5}$ under current $30 \mu\text{W cm}^{-2}$ illumination density. With light density increasing from 30 to $120 \mu\text{W cm}^{-2}$, the noise current increases from 1.18 to 2.80 nA $\text{Hz}^{-0.5}$ at -100 V bias.

Due to X-ray detection requires large voltage bias to collect charges induced in millimeter-thickness PSCs, here we characterize the noise current when the voltage bias varies

from -10 V to -200 V . Similar with **Figures 3B,D** shows the temporal response at 10 Hz with different voltages under $120 \mu\text{W cm}^{-2}$ illumination density at 460nm. The average noise current increases from 68.7 pA $\text{Hz}^{-0.5}$ to 3.24 nA $\text{Hz}^{-0.5}$ in darkness and increases from 0.25 nA $\text{Hz}^{-0.5}$ to 11.9 nA $\text{Hz}^{-0.5}$ under illumination as shown in **Figure 3E**. In conclusion, the function of light density and voltage vs. average noise is shown in **Figure 3F**. This device with epitaxial mono-crystalline films is an excellent candidate for X-ray detection in terms of their low signal noise performance under high voltage bias (Kasap et al., 2011).

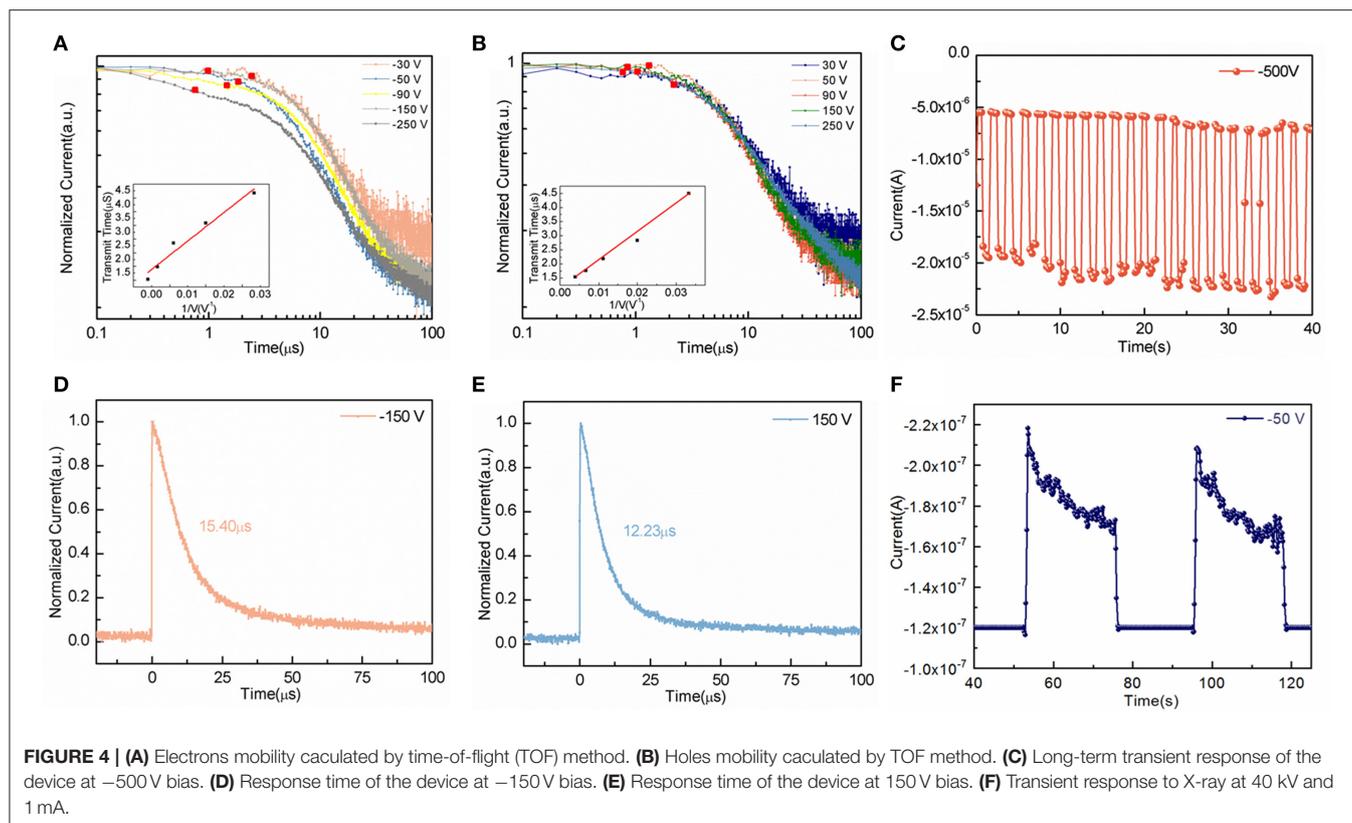
What's more, Responsivity (R), the ratio of photocurrent density and incident illumination density, indicates the efficiency of photodetectors responds to optical signal. And the formula is

$$R = \frac{J_{\text{light}} - J_{\text{dark}}}{P}$$

where J_{light} is photocurrent density, J_{dark} is dark current density and P is incident illumination density (Dou et al., 2014). The area of photodetector is 25 mm^2 . As illustrated in **Figure 3D**, R can be calculated to approximately 87 mA/W at -100 V and 23 mA/W at -50 V bias with $120 \mu\text{Wcm}^{-2}$ illumination at 460 nm.

Figures 4A,B provide the time of flight (TOF) curves to measure the carrier mobility of semiconductor material. According to

$$\mu = \frac{L^2}{T_{\text{tof}}V_s}$$



the carrier mobility μ can be directly calculated once we get the distance L and voltage bias V_s between two electrodes and the carrier transit time T_{tof} . (Huang et al., 2015; Wei et al., 2016) As depicted in **Figure 4A**, the average mobility of electrons is calculated to $385\text{ cm}^2\text{V}^{-1}\text{s}^{-1}$ under reverse voltages. Similarly, **Figure 4B** shows the average holes mobility of $415\text{ cm}^2\text{V}^{-1}\text{s}^{-1}$ under forward voltages. The mobility here is larger than the those of devices based on PSCs results from the decreased dislocation between epitaxial grown films and lattice matched substrate (Dong et al., 2015a; Ji et al., 2018).

Figure 4C illustrates stable dark current and photocurrent of the photodetector under high work voltage of -500 V bias at 460 nm . The long-term stability of the device benefits from epitaxial mono-crystalline films replacement of function layers (Wang et al., 2019). As shown in **Figures 4D,E**, the response time of device is $15.40\text{ }\mu\text{s}$ under reverse bias and $12.23\text{ }\mu\text{s}$ under forward bias, which is 3 orders of magnitude faster than those of photodetectors based on perovskite polycrystalline thin films (Liu et al., 2019; Su et al., 2019).

The temporal X-ray response of our device is illustrated in **Figure 4F**. With X-ray off, the dark current density is 480 nA cm^{-2} at -50 V bias. And the photocurrent density reaches 880 nA cm^{-2} to X-ray at 40 kV and 1 mA with dose rate of $6.7\text{ }\mu\text{Gys}^{-1}$. The sensitivity of our device is $59.7\text{ }\mu\text{C mGy}^{-1}\text{ cm}^{-2}$ which is almost five times higher than X-ray

detectors based on polycrystalline perovskite reported recently (Kim et al., 2017).

CONCLUSION

In summary, this work demonstrates a perovskite photodetector for X-ray detection fabricated with epitaxial mono-crystalline films as electron-blocking layers. Due to lattice matching, we fabricates the device with epitaxial layers of high crystallinity and considerable thickness, which are characterized by XRD and SEM. Furthermore, our X-ray photodetector shows low dark current density of $6.97 \times 10^{-11}\text{ A cm}^{-2}$ at -0.5 V bias, high electrons mobility of $385.9\text{ cm}^2\text{V}^{-1}\text{s}^{-1}$, high responsivity of 870 mA/W at -100 V bias and remarkable X-ray sensitivity of $59.7\text{ }\mu\text{C mGy}^{-1}\text{ cm}^{-2}$ at -50 V bias. This work offers a path of heteroepitaxial mono-crystalline films acting as electron-blocking layers for solution-processed photodetectors based on OIHPSCs.

DATA AVAILABILITY STATEMENT

The raw data supporting the conclusions of this article will be made available by the authors, without undue reservation.

AUTHOR CONTRIBUTIONS

YX, YL, and YP grew the perovskite single crystals. YX did the epitaxial experiments. YX and YP did the measurements. XW and YX analyzed these results. YX wrote this manuscript. All authors make comments on the manuscript.

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SUPPLEMENTARY MATERIAL

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Conflict of Interest: The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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