

# Glow-in-the-dark globular clusters: modelling their multiwavelength lanterns

C Venter<sup>1</sup>, I Buesching<sup>1</sup>, A Kopp<sup>2</sup>, A-C Clapson<sup>3</sup>, and O C de Jager<sup>1</sup>

<sup>1</sup>Centre for Space Research, North-West University, Potchefstroom Campus, Private Bag X6001, Potchefstroom 2520, South Africa

<sup>2</sup>Institut für Experimentelle und Angewandte Physik, Christian-Albrechts-Universität zu Kiel, Leibnizstrasse 11, 24118 Kiel, Germany

<sup>3</sup>Max-Planck-Institut für Kernphysik, PO Box 103980, 69029 Heidelberg, Germany

**Abstract.** Globular clusters (GCs) are astronomical tapestries embroidered with an abundance of exotic stellar-type objects, including ancient metal-poor stars, planetary nebulae, white dwarfs (WDs), low-mass X-ray binaries (LMXRBs), RR Lyrae variables, blue stragglers, cataclysmic variables, and possibly even central black holes. In addition, their high age promises a rich harvest of evolved stellar products, while the deep potential wells and high mass densities at their centres probably facilitate the formation of multiple-member stellar systems via increased stellar encounter rates. The ubiquity of GC LMXRBs, thought to be the progenitors of millisecond pulsars (MSPs), furthermore sets the stage for yet another interesting cluster subpopulation. In addition to the many GC radio pulsars and X-ray counterparts that have already been discovered, *Fermi* Large Area Telescope (LAT) recently unveiled the first gamma-ray GC pulsar (PSR J1823–3021A). The first observations of GCs in the GeV and TeV bands furthermore created much excitement, and in view of the above, it seems natural to explain these high-energy lanterns by investigating an MSP origin. An MSP population is expected to radiate several pulsed spectral components in the radio through gamma-ray wavebands, in addition to being sources of relativistic particles. The latter may interact with background photons in the clusters producing TeV excesses, while they may also radiate synchrotron photons as they traverse the cluster magnetic field. We will present our modelling results for the Terzan 5 cluster, focusing on the system constraints that may be derived in the context of this model by comparing our model to multiwavelength data. We also briefly discuss some alternative interpretations for the observed GC gamma-ray signals.

## 1. Introduction

There are about 150 Galactic globular clusters (GCs), ancient spherical arrangements of  $10^5 - 10^6$  stars bound by their mutual gravity. They prove useful in diverse astrophysical disciplines such as cosmology, galaxy formation, stellar evolution and dynamics, and binary and variable stars [24]. Empirical correlations between the number of low-mass X-ray binaries (LMXRBs), the progenitors of millisecond pulsars (MSPs), and the two-body stellar encounter rate [32], as well as the number of MSPs ( $N_{\text{MSP}}$ ) and this encounter rate [26, 2], provide evidence for a dynamical formation of such stellar objects in GCs. The high core densities of GCs are in turn believed to facilitate high encounter rates [42], leading to relatively large numbers of LMXBs and MSPs in GCs. These sources may be directly or indirectly responsible for, or contribute to, the multiwavelength emission seen from GCs. This paper discusses the multiwavelength observations and modelling of GCs, but with a strong focus on the gamma-ray waveband.

## 2. GCs by the dozen: *Fermi* Large Area Telescope (LAT) results

Just over a year into the *Fermi* mission, detection of 47 Tucanae at a  $17\sigma$  level was claimed, making this the first GC to be detected in gamma rays [1]. The spectrum was consistent with being an accumulation of several individual MSPs' pulsed spectra (for an alternative interpretation, see Section 6.2), with the measured luminosity constraining the 47 Tucanae MSP population to  $\sim 60$ . No pulsations from individual MSPs was however found. This was followed by the detection of Terzan 5 [27], as well as several (plausibly  $\sim 6$ ) other GCs [2], including M 28, NGC 6388, and Omega Cen. Another 3 detections and 3 GC candidates were claimed [35], bringing the total number of high-energy (HE) GCs to about a dozen.

## 3. The shy waveband: very-high-energy (VHE) observations

In contrast to the HE successes, TeV observations have only been able to uncover an excess in the direction of, but offset from the centre of, Terzan 5 [3]. The extended VHE source reaches beyond the cluster's tidal radius, exhibiting a power-law spectrum with a photon index of  $2.5 \pm 0.3_{\text{stat}} \pm 0.2_{\text{sys}}$  that implies an integral flux of  $(1.2 \pm 0.3) \times 10^{-12} \text{ cm}^{-2} \text{ s}^{-1}$  above 440 GeV. Only upper limits exist on other GC positions: 47 Tuc [7], M 13 [8, 29], M 5 and M 15 [29, 4], and NGC 6388 [4].

## 4. Multiwavelength pieces of the puzzle

GCs have now been detected in all energy domains [1]. There are various examples, including surface brightness profile measurements in the optical and infrared [36] which facilitate King-model-type fits that allow, e.g., inference of structural parameters; discoveries of embedded radio MSPs [34], X-ray MSPs, LMXBs, white dwarfs (WDs), and main-sequence binaries [21]; extended radio emission [15]; and the HE and VHE detections discussed above.

Diffuse X-rays have been reported for M 22, Omega Cen, 47 Tucanae [25, 28], M 80, NGC 6266, NGC 6752, and M 5 [30]. This emission has been interpreted as being due to the formation of bow shocks which result from the GC motion through the Galactic halo plasma, since it appears to be associated with GCs with a high proper motion and large accumulation of internal gas. However, the clumpy structure observed near NGC 6752 exhibits a hard non-thermal spectrum as well as a radio counterpart, and may be due to bremsstrahlung by shock-accelerated electrons hitting nearby gas clouds, while unresolved X-ray sources may contribute or even dominate such radiation in general [17].

Diffuse X-ray emission has also been observed from Terzan 5 in the 1–7 keV band [17]. This radiation peaked at the cluster centre and decreased smoothly outwards. Observational results favoured a non-thermal origin of these X-rays (e.g., synchrotron radiation – SR). A follow-up search for diffuse X-rays from six HE GCs yielded no significant emission above the background level [18]. Within the MSP scenario that this paper focuses on, SR is expected to be produced by relativistic leptons that escape from a population of MSPs inside GCs and interact with the GC magnetic field [37]. Using the diffuse X-ray profile [17], the diffusion coefficient of these particles may be constrained [14]. One can also probe the cluster magnetic field's magnitude and profile.

## 5. Message from the most powerful gamma-ray MSP

The recent detection [19] by *Fermi* of the most powerful (and likely youngest) gamma-ray MSP to date, also being the first firm<sup>1</sup> detection of a gamma-ray MSP in a GC, allows new constraints on the underlying mechanism responsible for the observed HE GC emission (see Section 2). Using the measured gamma-ray light curve, one may define an ‘off-pulse window’

<sup>1</sup> *AGILE* has claimed a  $4\sigma$  detection of PSR J1824–2452 in the GC M 28 [31], but to date no such pulsations have been confirmed by *Fermi* [2, 19].

– a range in the cyclic normalized phase coordinate where the pulse is believed to be ‘switched off’ ( $0.07 < \phi < 0.60$  and  $0.67 < \phi < 0.90$  in this case). Selecting off-pulse gamma rays only, one can then constrain the level of emission from other sources (e.g., other GC MSPs in this instance) or from the background. This technique is also used when searching for steady emission from putative pulsar wind nebulae surrounding younger pulsars [5]. *Fermi* did not detect any GeV point sources in the off-pulse region of PSR J1823–3021A. A previous estimate [35, 2] assigned  $\sim 100$  gamma-ray MSPs to this cluster, assuming typical MSP spectral properties. We now know that this extraordinarily powerful MSP accounts for almost all of the HE emission coming from NGC 6624, and the off-pulse flux upper limit constrains the number of additional ‘ordinary’ MSPs hosted by this cluster to  $< 32$ . This implies that, at least for NGC 6624, the HE emission must be predominantly due to pulsed emission, believed to be magnetospheric curvature radiation (CR; Section 6.1). Any unpulsed emission (e.g., due to inverse Compton (IC) scattering) is severely constrained in this case (see Section 6.2). Although this supports a pulsed origin for HE emission in other GCs, one can only speculate as to the general validity, given the plausible differences in MSP population properties and environments between GCs (i.e., the relative contribution of pulsed and unpulsed emission is generally unknown).

## 6. Successes and challenges of the gamma-ray MSP population model

### 6.1. Pulsed HE emission

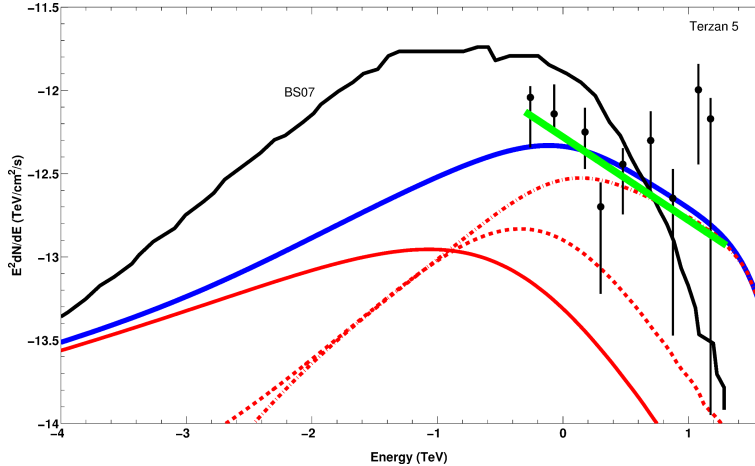
The fact that *Fermi* has detected several GCs exhibiting spectra that are very reminiscent of pulsar spectra (hard power law with exponential cutoff at a few GeV) leads to the notion that a population of MSPs hosted within the GC may be cumulatively responsible for this HE emission [2], strengthened by the discovery of gamma rays from PSR J1823–3021A (see Section 5). Indeed, this has been predicted [23, 38, 39] shortly before the discovery of 47 Tucanae in the HE band [1]. The pulsed emission arises from particles being accelerated by electric fields inside the MSP magnetospheres, and suffering CR losses as they move along curved magnetic field lines, before being ejected beyond the light cylinder. Predictions were reasonably close to the measured spectrum, and implied a population size of  $\sim 50$  members [41], but overpredicted the spectral cutoff by a factor  $\sim 2$ , suggesting some revision of the model (e.g., electric field and curvature radius). The predicted spectrum was also too hard [39]. It has to be borne in mind, however, that a pair-starved electric field has been assumed in these calculations. Subsequent light curve modelling of several gamma-ray MSPs [40] suggests that the bulk of this population may have screened electric fields, a possibility not considered previously within the standard assumption of MSP magnetospheres consisting of dipolar magnetic fields [22].

### 6.2. Unpulsed HE emission

As an alternative to the CR mechanism, an IC scenario was considered to explain the HE fluxes seen by *Fermi* [14]. The model solves a cosmic-ray diffusion equation (using a slightly different stellar photon energy density profile than [10]), and predicted that the bulk of the HE radiation comes from a region beyond the GC core (contrary to the findings of [10]), so that GCs should be extended HE sources. The *Fermi* flux may be reproduced for some combinations of model parameters in this scenario. This model also predicts spectral components that should be visible in the VHE domain in some cases. Such unpulsed IC components seem less dominant, however, in the case of the cluster NGC 6624, as explained in Section 5.

### 6.3. Unpulsed VHE emission

In addition to being sources of pulsed photons, the large potential drops in MSP magnetospheres indicate that they are also sources of relativistic leptons: MSPs may produce leptons with energies of up to a few TeV [13]. These leptons escape from the magnetospheres, and may be reaccelerated in shocks formed by collisions of stellar winds in the cluster core [10, 11]. The



**Figure 1.** Predicted IC spectrum in the MSP scenario. The black solid line labelled BS07 is scaled from [10], while the thick blue solid line represents the sum of the IC emission from three zones lying at increasing radii from the GC centre (indicated by the lower solid, dashed, and dashed-dotted red lines). We assumed a power-law lepton injection spectrum with minimum energy of 100 MeV, maximum energy of 100 TeV, photon index of 1.6, cluster field of  $1 \mu\text{G}$ , and total power of  $2 \times 10^{34} \text{ erg s}^{-1}$  (e.g., a total number of 100 MSPs if the conversion efficiency of spin-down power to particle power is 1%). We also introduced a factor 3 scaling. This implies a larger value of  $N_{\text{MSP}}$ , number of cluster stars ( $N_{\text{star}} > 8 \times 10^5$ ), average stellar radius ( $R > 10^6 \text{ cm}$ ), stellar temperature ( $T > 6000 \text{ K}$ ), or smaller GC distance ( $d < 5.9 \text{ kpc}$ ), or a combination of these [33]. The green line and data points are from H.E.S.S. results [3].

leptons may then upscatter bright starlight and cosmic microwave background (CMB) photons to very high energies, leading to an unpulsed VHE spectral component, which gives an independent constraint<sup>2</sup> (vs. pulsed emission) on  $N_{\text{MSP}}$  (depending on the cluster magnetic field and diffusion coefficient) [39]. It was found that GCs such as 47 Tucanae and Terzan 5 may be visible for H.E.S.S., depending on model parameters [10, 39]. Using the H.E.S.S. upper limits on the VHE gamma-ray emission from 47 Tucanae, we could infer a population of  $\sim 30 - 40$  MSPs, given a cluster magnetic field of  $B \sim 10 \mu\text{G}$  (but quite larger for  $B < 5 \mu\text{G}$  or  $B > 30 \mu\text{G}$ ).

We have now extended our calculations for Terzan 5 (Figure 1), including a third large emission zone extending up to the tidal radius, and a full calculation of the energy density profile [33], and assuming a power-law injection spectrum for the leptons. It has been expected that, since the stellar energy density profile is strongly peaked at the centre of the GC, the TeV flux should follow a similar profile, as the IC process depends on this target soft photon field [16]. However, even though the energy density drops steeply as one leaves the core region, the much larger size of the halo traps the particles for much longer, where they interact with the low-energy-density field. This increase in residence time outweighs the drop in soft photon energy density, and emission from this region dominates the VHE spectrum in our model (dashed-dotted line in Figure 1).

#### 6.4. Challenges

Typical constraints derived from the measured energy fluxes in the GeV band on  $N_{\text{MSP}}$  contained within a GC are not always so restrictive, given uncertainties in energy flux, GC distance, and average beaming factors. This means that the level of the pulsed CR flux predictions is

<sup>2</sup> We are assuming that the number of visible HE MSPs  $N_{\text{vis}} \approx N_{\text{MSP}}$ .

somewhat uncertain. Lack of evidence for spectral cutoffs around a few GeV for some GCs also challenge the cumulative CR interpretation for the HE GC emission in these cases [35, 12]. In the case of the VHE models, refinement of the soft photon energy density profile (e.g., adding a Galactic component [14]), cluster magnetic field profile, and particle transport calculations may be needed. The peculiar asymmetric, offset VHE source seen in coincidence with Terzan 5 also challenges a simplistic MSP scenario. However, such an offset may be introduced by several factors, including a small MSP population at relatively large distances from the centre, formation of MSPs near the tidal radius in addition to the core [35], non-spherical photon target fields, proper motion of the GC, and an asymmetric diffusion coefficient [14]. Source extension, morphology, spectrum and energetics will continue to constrain GC models.

## 7. Other sides of the VHE coin: alternative explanations

### 7.1. Explosive energy: a short gamma-ray burst (GRB) relic in Terzan 5?

Long GRBs ( $> 2$  s), signalling the death of massive stars [20], have been invoked as a possible origin for some of the many unidentified TeV sources observed by H.E.S.S. [9, 6]. Similarly, the scenario of the VHE emission from Terzan 5 being due to the remnant of a short ( $< 2$  s), powerful GRB (thought to be due to a compact binary merger [20]) has been put forward [16]. The main argument relies on the fact that the high-density environments of ancient GCs are conducive to the formation of compact binaries, and that the merger of the members (e.g., two neutron stars) of some may accelerate hadronic cosmic rays. The latter may interact with ambient target nuclei and decay to gamma rays via the  $\pi^0$  process. It is shown that for a target density of  $n \sim 0.1 \text{ cm}^{-3}$ , and a broken power-law cosmic-ray spectral shape with spectral index of 2.0 below 5 TeV, the total energy of hadrons of  $E \approx 10^{51}$  ergs may possibly be supplied by ultrarelativistic blast waves converting a significant part of the kinetic energy to cosmic ray particle energy. A break in the gamma-ray spectrum in the GeV / TeV range would support such a scenario. The putative remnant age of  $\sim 10^4$  yr (estimated from the TeV source extension and depending on the hadron diffusion coefficient) corresponds roughly to the estimated rate of short bursts in the Milky Way ( $\sim 1$  event per  $10^4$  yr, depending on the burst beaming factor) in the event of relatively slow diffusion. In this scenario, electrons accelerated by the blast wave scatter the stellar photons to produce diffuse non-thermal X-ray emission. Additional predicted signatures are faint thermal X-rays from hot thermal plasma in the remnant, faint nuclear line emission from radio-active decay of heavy nuclei ejected during the merger, and faint ionization lines following from interaction of the GRB afterglow with the interstellar medium.

### 7.2. Distant cousins: a partly WD origin?

Instead of MSPs, a population of fast-rotating, magnetized WDs contained within the GC has been investigated as being responsible for (part of) the HE and VHE signals observed from these clusters [12]. An estimate involving the initial mass function implies that the number of all WDs in a GC may be as high as  $\sim 10^5$ , dominating the MSP population by a factor of a few hundred. The evolutionary behaviour of the non-accreting WD energetics is furthermore very similar to the case of rotation-powered pulsars, although the typical surface magnetic fields may be lower ( $\sim 10^8 \text{ G} < 10^{8-9} \text{ G}$ ), rotational periods much longer ( $\sim 100 \text{ s} > 5 \text{ ms}$ ), masses similar ( $\sim 0.8M_{\odot} \approx 1.4M_{\odot}$ ), and radii larger ( $\sim 5 \times 10^8 \text{ cm} > 10^6 \text{ cm}$ ). These parameters suggest that a single WD may have injected electrons with a power of up to  $\sim 10^{28} \text{ erg s}^{-1}$  (vs.  $\sim 10^{32} \text{ erg s}^{-1}$  for an MSP), or less in the case of magnetic field decay. In addition, WDs created by WD-WD mergers in compact binary systems (having larger masses, smaller radii, and shorter periods than non-accreting WDs) may also contribute to the observed gamma-ray flux, but at a much lower level (a factor  $\sim 10$ ). It has been shown [12] that a few thousand WDs that have been created uniformly during the GC lifetime may produce a detectable VHE signal via upscattering of CMB and stellar soft photon fields (assuming mono-energetic electron injection spectra, and

different models for WD formation, evolution, and magnetic field decay, and also depending on the electron acceleration process).

## 8. Conclusion

We have discussed GCs as multiwavelength objects, focusing on their gamma-ray properties. Although some models prove reasonable, inevitable refinement is due in the wake of observations increasing both in quality and quantity.

## Acknowledgments

This research is based upon work supported by the South African National Research Foundation.

## References

- [1] Abdo A A *et al.* 2009 *Science* **325** 845–8
- [2] Abdo A A *et al.* 2010 *Astron. & Astrophys.* **524** A75
- [3] Abramowski A *et al.* 2011 *Astron. Astrophys.* **531** L18–22
- [4] Abramowski A *et al.* 2011 *Astrophys. J.* **735** 12–9
- [5] Ackermann M *et al.* 2011 *Astrophys. J.* **726** 35–51
- [6] Aharonian F *et al.* 2008 *Astron. & Astrophys.* **477** 353–63
- [7] Aharonian F *et al.* 2009 *Astron. Astrophys* **499** 273–7
- [8] Anderhub H *et al.* 2009 *Astrophys. J.* **702** 266–9
- [9] Atoyan A, Buckley J and Krawczynski H 2006 *Astrophys. J.* **642** L153–6
- [10] Bednarek W and Sitarek J 2007 *Mon. Not. Royal Astron. Soc.* **377** 920–30
- [11] Bednarek W 2011 *High-Energy Emission from Pulsars and their Systems: Proc. 1st Session Sant Cugat Forum of Astrophys.* ed N Rea and D F Torres pp 185–205 arXiv:1009.1694
- [12] Bednarek W 2012 *J. Phys. G: Nucl. Part. Phys.* **39** 065001
- [13] Büsching I, Venter C and de Jager O C 2008 *Adv. Space Res.* **42** 497–503
- [14] Cheng K S *et al.* 2010 *ApJ* **723** 1219–30
- [15] Clapson A-C, Domainko W F, Jamrozy M, Dyrda M and Eger P 2011 *Astron. Astrophys.* **532** A47
- [16] Domainko W F 2011 *Astron. Astrophys.* **533** L5–8
- [17] Eger P, Domainko, W F and Clapson, A-C 2010 *Astron. Astrophys.* **513** A66
- [18] Eger P and Domainko W F 2012 *Astron. Astrophys.* **540** A17
- [19] Freire P C C *et al.* 2011 *Science* **334** 1107–10
- [20] Gehrels N *et al.* 2005 *Nature* **437** 851–4
- [21] Grindlay J E, Heinke C, Edmonds P D and Murray S S 2001 *Science* **292** 2290–5
- [22] Harding A K, Muslimov A G and Zhang B 2002 *Astrophys. J.* **576** 366–75
- [23] Harding A K, Usov V V and Muslimov A G 2005 *Astrophys. J.* **622** 531–43
- [24] Harris W E 1996 *Astron. J.* **112** 1487–8
- [25] Hartwick F D A, Grindlay J E and Cowley A P 1982 *Astrophys. J.* **254** L11–3
- [26] Hui C Y, Cheng K S and Taam R E 2010 *Astrophys. J.* **714** 1149–54
- [27] Kong A K H, Hui C Y and Cheng K S 2010 *Astrophys. J.* **71** L36–9
- [28] Krockenberger M and Grindlay J E 1995 *Astrophys. J.* **451** 200–9
- [29] McCutcheon M *et al.* 2009 *Proc. 31st ICRC, Lodz, Poland* arXiv:0907.4974
- [30] Okada Y, Kokubun M, Yuasa T and Makishima K 2007 *Publ. Astron. Soc. Japan* **59** 727–42
- [31] Pellizzoni A *et al.* 2009 *Astrophys. J.* **695** L115–9
- [32] Pooley D *et al.* *Astrophys. J.* **591** L131–4
- [33] Prinsloo L P, Venter C, Buesching I and Kopp A *these proceedings*
- [34] Ransom S M 2008 *40 Years of Pulsars: Millisecond Pulsars, Magnetars and More, AIP Conf. Ser., ed C Bassa, Z Wang, A Cumming, V M Kaspi* **983** pp 415–23
- [35] Tam P H T *et al.* 2011 *Astrophys. J.* **729** 90–7
- [36] Trager S C, King I R and Djorgovski S 1995 *Astron. J.* **109** 218–41
- [37] Venter C and de Jager O C 2008 *AIP Conf. Ser.* **1085** 277–80
- [38] Venter C and de Jager O C 2008, *Astrophys. J.* **680** L125–8
- [39] Venter C *et al.* 2009 *Astrophys. J.* **696** L52–5
- [40] Venter C *et al.* 2009 *Astrophys. J.* **707** 800–22
- [41] Venter C, de Jager O C, Kopp A and Büsching I 2011 *2011 Fermi Symp. Proc. eConf C110509* arXiv:1111.1289
- [42] Verbunt F and Hut P 1987 *IAU Symposium ed D J Helfand and J-H Huang* **125** 187–96